

Existence of Critical Gompertz Parameter for Solid Tumour Growth Model and Its Asymptotic Expression

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Abstract—In this paper we provide an interval of existence of critical Gompertz parameter of solid tumour growth model and their asymptotic formula for large number of tumour cells, in the absence of specific volume data at particular time.

I. INTRODUCTION

Cancer is a disease that may affect people at all ages. It causes about 13% of all human deaths. In the past two decades increasing attention has been paid to tumour growth, a major death cause in our society. It has long been documented that the growth of cancer tumours follows a sigmoidal growth curve, exhibiting at first a phase of exponential growth and later a phase of slowed growth. The character of such curves is simulated by various mathematical formulations, including deterministic models, stochastic models and cellular automata models. The most recognized pattern of tumour growth is Gompertz growth, which has been utilized by many researchers to provide a basis of description and prediction [1-17].

Broadly speaking solid tumour growth may be termed avascular or vascular with angiogenesis facilitating the transformation from avascular to vascular growth [18-21]. The avascular stage can be characterised by diffusion-limited growth, with the tumour receiving vital nutrients and eliminating waste products via diffusion across its outer boundary [21-23]. Since the size to which such tumours grow is limited ($O(\text{mm}^3)$), avascular tumours are usually harmless. To escape from the restrictions of avascular growth a tumour must undergo angiogenesis [18-21]. During this process the tumour induces blood vessels from the surrounding tissue to form a new capillary network that migrates towards, and ultimately penetrates, the tumour. Once vascularised the tumour has access to an almost limitless supply of nutrients and is potentially life-threatening for two reasons. Firstly, the rapid growth that results may impair the function of vital organs. Second, the development of secondary tumours or metastases is now a real threat; tumour fragments that enter the blood supply are transported to other parts of the body where if conditions are favourable they establish secondary tumour colonies that further threaten the host.

Traditionally mathematical models describing avascular tumour growth assume radial symmetry of the tumour and focus on its responses to various growth factors [24-26]. These models show excellent agreement with experimental results, reproducing the multi-layered structures that characterise avascular tumours and multicellular spheroids.

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However, the deterministic Gompertz law of population growth has been widely used to describe in vivo tumour growth in experimental oncology [27-32]. If $V(t)$ is the volume of the tumour cell at time t , then the Gompertz law models the cell growth by the equation

$$\frac{dV(t)}{dt} = AV(t) - \beta V(t) \ln V^*(t) \quad (1)$$

where A ; the intrinsic growth rate of the tumor, is a parameter related to the initial mitosis rate and β ; the growth deceleration factor and $V^*(t) = V(t)/V_0$; define $V(0) = V_0$, is the volume at time $t = 0$. From a biological point of view, a greater β value or a smaller A value indicates a greater antitumoral effect of the therapy [33].

The plan of this paper is as follows. In section 2 the details of a procedure for the estimation of parameters A and β in the absence of specific volume data. This method utilizes the cumulative volume rate (V_c) (i.e., defined by $V_c = \int_0^{\infty} V^*(t) dt$, where $V^*(t) = V(t)/V_0$

and the maximum lifetime of tumour cells (t_m) (at this time the tumour reaches its maximum size of volume or maximum number of cells before disintegration and final effects). Also given details about behaviour growth critical Gompertz parameter of solid tumour and proved the existence of critical time t_k , critical Volume N_k . Section 3 presents the results for interval of existence of numerical solution. In section 4 we given a conclusion.

II. ESTIMATION OF PARAMETERS

An exact mathematical description of our model of tumour cell proliferation is given by a Gompertz equation (1) of the following form

$$V(t) = V_0 e^{\frac{A}{\beta}(1-e^{-\beta t})} \quad (2)$$

where $V(t)$ is the clonogenic tumour volume at time t ; V_0 is the clonogen number at time $t = 0$: A and $\beta (> 0)$ are the Gompertz growth parameters.

The doubling time is a key parameter for assessing the impact of delays in cancer treatment. Most of the information about tumour growth rates comes from studies performed long ago and not known clearly the maximum volume size of individual tumours and groups of tumours. In general the time the tumour takes to double itself varies widely, such that in case of histological type of tumour the time distribution for tumour doubling itself is normally long [34-38]. The Gompertz model presents a doubling time (Volume Rate Doubling time (VRD)) which depends only on β .

Comparisons of volume data of solid tumours in tumour growth model are aided by calculation



of the VRD, because VRD changes in the same direction as lifespan of tumour cells. The growth rate of the tumour may also be described by additional coefficients (Gompertz-Makeham model) or by other power functions (Weibull model), in which the VRD changes with time [37]. Solving equation(2) for VRD gives

$$VRD = -\frac{1}{\beta} \ln \left[1 - \frac{\beta}{A} \ln(2) \right] \tag{3}$$

The above equation (3) fully depends on β ; so we have to estimate β to calculate VRD in the absence of specific volume data of solid tumour cells, since VRD changes in the same direction as lifespan of tumour cells. Equation (2) gives

$$\frac{A}{\beta} = \frac{\ln(V^*(t))}{(1 - e^{-\beta t})} \tag{4}$$

where $V^*(t) = V(t)/V_0$. We assume that V (tm); maximum volume of tumour cells (where tm is the time at which the tumour contains a cell number which is one less than its maximum i.e., one cell less to death, and which approximates the maximum lifespan of tumour cells $t^*(m)$). Thus

$$V(t^*_m) = V_0 e^{\frac{A}{\beta}(1 - e^{-\beta t^*_m})} \tag{5}$$

after a few algebraic manipulations we get,

$$t^*_m = -\frac{1}{\beta} \ln \left[1 - \frac{\beta}{A} \ln \left(\frac{V(t^*_m)}{V_0} \right) \right] \tag{6}$$

The cumulative intrinsic volume growth rate V_c of the Gompertz model of equation (2), is

defined $V_c = \int_0^\infty V^*(t) dt$ Substitute the value of

$V^*(t)$ from the equation (2) in the above equation and apply a little algebra we get the following equation

$$-\beta = \frac{1}{V_c} e^{-\frac{A}{\beta}} \int_{\frac{A}{\beta}}^{\infty} \frac{e^{-z}}{z} dz \tag{7}$$

Where $z = -\frac{A}{\beta} e^{-\beta t}$. Clearly, the above integral (7),

exists $\forall \beta \in R$.

Consider the initial volume of size V_0 at initial stage. From the equation (1) the volume at which the growth rate of initial volume V_0 has increasing or, equivalently, it tends to a V ($t^*(m)$) is the time at which volume approximates the maximum number of tumour cells, is called a critical time, t_k . The remaining tumour cells from an original volume size V (t) surviving at this critical time is called critical volume V_k , and the corresponding Gompertz parameter in equation (1), is called critical Gompertz parameter β_k .

2.1 Behaviour of solid tumour growth Gompertz parameter To find the critical points of β we consider the partials of β with respect to $V^*(t), V_c$ and t_m in equation (7). These are given by

$$\frac{\partial \beta}{\partial V^*(t)} = \frac{\left[A - \left(\frac{1}{V_c} \right) \right] / V^*(t) \ln V^*(t)}{1 + \left(\frac{e^{-\beta t_m}}{(e^{-\beta t_m} - 1)} \right) t_m \left[A - \left(\frac{1}{V_c} \right) \right]} \tag{8}$$

$$\frac{\partial \beta}{\partial V_c} = \frac{-\beta / V_c}{1 + \left(\frac{e^{-\beta t_m}}{(e^{-\beta t_m} - 1)} \right) t_m \left[A - \left(\frac{1}{V_c} \right) \right]} \tag{9}$$

and

$$\frac{\partial \beta}{\partial t_m} = \frac{-\beta \left(\frac{e^{-\beta t_m}}{(e^{-\beta t_m} - 1)} \right) / \left[\left(\frac{1}{V_c} \right) - A \right]}{1 + \left(\frac{e^{-\beta t_m}}{(e^{-\beta t_m} - 1)} \right) t_m \left[A - \left(\frac{1}{V_c} \right) \right]} \tag{10}$$

Here, $\beta t_m \geq 1$ and $[A \leq (1/V_c)]$ is positive, since $\left(\frac{1}{V_c} \right) \leq A$

Therefore, for $\beta t_m \geq 1$ the value of

$1 + \left(\frac{e^{-\beta t_m}}{(e^{-\beta t_m} - 1)} \right) t_m \left[A - \left(\frac{1}{V_c} \right) \right]$ is positive.

This will give the result

$$\frac{\partial \beta}{\partial V^*(t)} \geq 0, \frac{\partial \beta}{\partial V_c} \leq 0 \text{ and } \frac{\partial \beta}{\partial t_m} \leq 0$$

It shows that the value of parameter β is increases when value of $V^*(t)$ is increases, β is decrease when the value of V_c increase and also β is decrease when the value of t_m increases.

If we send $V^*(t)$ to ∞

in (8),(9) and (10) get that

$$\lim_{V^*(t) \rightarrow \infty} \frac{\partial \beta}{\partial V^*(t)} = 0, \lim_{V^*(t) \rightarrow \infty} \frac{\partial \beta}{\partial V_c} = 0 \text{ and } \lim_{V^*(t) \rightarrow \infty} \frac{\partial \beta}{\partial t_m} = 0.$$

Also we obtain,

$$\frac{\partial \beta}{\partial V^*(t)} = 0 \Leftrightarrow A = \frac{1}{V_c}, \quad \forall V^*(t),$$

$$\frac{\partial \beta}{\partial V_c} = 0 \Leftrightarrow \beta = 0,$$

$$\frac{\partial \beta}{\partial t_m} = 0 \Leftrightarrow A = \frac{1}{V_c} \text{ or } \beta = 0 \text{ or } \beta = -\infty.$$

Upon substitution $A = \frac{1}{V_c}$ in (4), we get

$$\frac{t_m}{V^*(t) \ln V^*(t)} = \frac{\beta t_m}{(1 - e^{-\beta t})} \tag{11}$$

from (11) it follows that,

$$\frac{t_m}{V^*(t) \ln V^*(t)} \begin{cases} < 1 \text{ if } \beta < 0 \\ = 1 \text{ if } \beta = 0 \\ > 1 \text{ if } \beta > 0 \end{cases} \tag{12}$$

The point

$$\frac{t_m}{V_c \ln V^*(t)} = 1$$

at which β changes sign is said to be the critical point β_k .

From [39], Collins and his co-workers [40] were able to show that for a series of 206 children with Wilms' tumour that risk of recurrence agreed well with theoretical prediction by the method of Boag [41] and also the growth rate function approaches a constant with predictions on the basis of exponential growth at larger time. To determine the critical

Gompertz parameter, β_k , first we use the identity $\frac{t_m}{V_c \ln V^*(t)} = 1$ to obtain critical values of β ; namely critical volume, V_k and critical time t_k . Since the partials of β with respect to

$V^*(t), \frac{1}{V_c}$ and t_m become zero at $\frac{t_m}{V_c \ln V^*(t)} = 1$.

Note that the condition $\frac{t_m}{V_c \ln V^*(t)} = 1$

is necessary to have a constant growth rate function. Finally we obtain the asymptotic solution of (7), for the critical values, V_k and t_k . Now we shall prove the existence of critical volume and critical time.

2.2 Critical volume V_k

For a given

$V^*(t), \frac{1}{V_c}$ and t_m with

$$\frac{t_m}{V_c \ln V^*(t)} > 1$$

There exists a critical volume V_k and is given by

$e^{\frac{t_m}{V_c}}$. Indeed, since $\ln V^*(t) < \frac{t_m}{V_c}$ we can take $\ln V^*(t) = \frac{t_m}{V_c}$, or

$V_k = e^{\frac{t_m}{V_c}}$. For instance,(see Table

I in [42] when (Mouse Krebs)) $A = 5:25, t_m = 15; 20$ and 25 , we find that $V_k = 927 \times 10^6, 949 \times 10^6$ and 952×10^6 , respectively. Note that V_k is increases as t_m . Thus the remaining volume (critical

volume)approximates 949×10^6 to 952×10^6 cells. Since the volume reaches its maximum size, the above said t_m can be treated as critical time t_k . From this we obtained $t_k = V_c \ln [V_k]$. To study the tumor growth rate of the remaining critical volume we need to consider critical Gompertz parameter β_k because

$$\frac{t_m}{V_c \ln V^*(t)} = 1$$

when $V^*(t) = V_k$. Thus, we conclude that when

$$\frac{t_m}{V_c \ln V^*(t)} > 1, t_k = t_m \text{ and } V_k = e^{\frac{t_m}{V_c}}. \text{ Clearly, when } \frac{t_m}{V_c \ln V^*(t)} > 1, \text{ both } t_k \text{ and } t_m \text{ are same.}$$

2.3 Critical growth time t_k

On the contrary, when

$$\frac{t_m}{V_c \ln V^*(t)} < 1$$

it is trivial to find the critical volume. As $t_m < V_c \ln [V_k]$ we can take $t_k = V_c \ln [V_k]$. For instance,(see Table I in [42] when (Rat R39 Sarcoma,R3a7))

$A = 1.28, t_m = 42.44$ days and $V^*(t) = 241 \text{ cm}^3$, we find that $t_k = 28.4854$. Thus we conclude that when

$$\frac{t_m}{V_c \ln V^*(t)} < 1$$

., $t_k = V_c \ln [V_k], V_k = e^{\frac{t_m}{V_c}}$ and $t_k \neq t_m$.

Clearly, when $\frac{t_m}{V_c \ln V^*(t)} < 1, t_k < t_m$. Thus, in general, for any given $\frac{1}{V_c}, t_m$ and $V^*(t)$, we get $t_k \leq t_m$.

III. INTERVAL OF EXISTENCE OF NUMERICAL SOLUTION

Recall that in [42-44] we obtained asymptotic solution of equation(7) for a large $V^*(t)$ with $\frac{t_m}{V_c \ln V^*(t)} > 1$ for $\beta > 0$. The asymptotic solution is given by

$$\beta = -\frac{1}{t_m} \ln \left[1 - e^{\left(\frac{t_m}{V_c}\right)C / \left(\frac{t_m}{V_c}-1\right)} (\ln V^*(t))^{\left(\frac{t_m}{V_c}\right) / \left(\frac{t_m}{V_c}-1\right)} \right] \quad (13)$$

and

$$\beta = -\frac{1}{t_m} \ln \left[1 - \frac{C \ln V^*(t) - \ln V^*(t) - 1}{\ln(\ln V^*(t)) + C} \right], \quad (14)$$

when $\frac{t_m}{V_c} \neq 1$ and $\frac{t_m}{V_c} = 1$ respectively, and

$C = 0.577215$; Euler's constant.

We remark that when $\frac{t_m}{V_c \ln V^*(t)} < 1$

the above asymptotic formulae remain valid, provided $\frac{t_m}{V^*(t) \ln V^*(t)} < 1$ or $V^*(t) > V_k$.

Our immediate concern is to extend the above formulae to $V^*(t) \geq V_k$. It is amazing to learn that the formulae (13) and (14) derived for a large $V^*(t)$ equally hold good for $V^*(t) = V_k$. Here is the proof: Upon substitution $V^*(t) = V_k$ to (7), we get

$$\beta_k = \frac{-1}{V_c} e^{\frac{-\ln(V_k)}{(1-e^{-\beta_k t_k})}} \int_{-\ln(V_k)}^{\infty} \frac{e^{-z}}{z} dz \quad (15)$$

where $t_k = t_m, \frac{t_m}{V_c \ln V^*(t)} < 1$

and $t_k = V_c \ln [V_k]$ when $\frac{t_m}{V_c \ln V^*(t)} < 1$

Equation (15) gives

$$\begin{aligned} \beta_k &\leq \frac{-1}{V_c} e^{\frac{-\ln(V_k)}{(1-e^{-\beta_k t_k})}} \left(\frac{-(1-e^{-\beta_k t_k})}{\ln(V_k)} \right) \int_{\frac{-\ln(V_k)}{(1-e^{-\beta_k t_k})}}^{\infty} \frac{e^{-z}}{z} dz \\ &= \frac{-1}{V_c} e^{\frac{-\ln(V_k)}{(1-e^{-\beta_k t_k})}} \left(\frac{-(1-e^{-\beta_k t_k})}{\ln(V_k)} \right) e^{\frac{\ln(V_k)}{(1-e^{-\beta_k t_k})}} \end{aligned}$$

Hence

$$\beta_k t_k \leq \frac{t_k}{V_c} \left(\frac{1-e^{-\beta_k t_k}}{\ln(V_k)} \right).$$

The above inequality ensures the existence of solution β_k , provided

$$\frac{t_m}{V_c \ln V^*(t)} = 1 \text{ As a consequence, we get } \beta_k t_k \leq (1 - e^{-\beta_k t_k}). \quad (16)$$

Obviously, any $\beta_k \geq 0$ will satisfy (16). To obtain the solution of (15) for every fixed V_c, t_k and V_k choose any β^* from the solution interval $[0, \infty)$. As every $\beta_k \geq \beta^*$ satisfies (15), let us find the largest possible of these. To achieve this, fix one such β^* . Substitute this into the right hand side of (15) to get

$$\beta_k = \frac{-1}{V_c} e^{\frac{-\ln(V_k)}{(1-e^{-\beta^* t_k})}} \int_{-\ln(V_k)}^{\infty} \frac{e^{-z}}{z} dz \quad (17)$$

where

$\epsilon = (1 - e^{-\beta^* t_k}) / (1 - e^{-\beta_k t_k}) \leq 1$, since $\beta_k \geq \beta^*$. We observe that by sending ϵ to zero in (17), one can obtain the asymptotic value of β_k for a large $V^*(t) = V_k$, as we have

$$0 \leq \beta_k t_k \leq \frac{t_k}{V_c} \left(\frac{1 - e^{-\beta^* t_k}}{\epsilon} \right), \quad (18)$$

from(16).

If we send ϵ to zero, on account of 8.212,1. in [45], we get

$$\beta_k = \lim_{\epsilon \rightarrow 0} -\frac{1}{t_k} \ln \left[1 - \left(\frac{1 - e^{-\beta^* t_k}}{\epsilon} \right) \right] \quad (19a)$$

$$= \lim_{\epsilon \rightarrow 0} -\frac{1}{t_k} \ln \left[1 - \left(\frac{1 - e^{-\beta^* t_k}}{\epsilon} \right) \right] - \lim_{\epsilon \rightarrow 0} -\frac{1}{t_k} \ln \left[\left(\frac{\epsilon}{1 - e^{-\beta^* t_k}} \right) - 1 \right] \quad (19b)$$

And also from equation (17),

$$\beta_k = \lim_{\epsilon \rightarrow 0} \frac{-1}{V_c} e^{\left(\frac{-\ln(V_k)}{(1-e^{-\beta^* t_k})}\right)\epsilon} \int_{\left(\frac{-\ln(V_k)}{(1-e^{-\beta^* t_k})}\right)\epsilon}^{\infty} \frac{e^{-z}}{z} dz \quad (20a)$$

$$= \lim_{\epsilon \rightarrow 0} \frac{-1}{V_c} e^{\left(\frac{-\ln(V_k)}{(1-e^{-\beta^* t_k})}\right)\epsilon} \left[C + \ln \left[\frac{\ln(V_k)}{1 - e^{-\beta^* t_k}} \right] \epsilon + \int_0^{\left(\frac{\ln(V_k)}{1 - e^{-\beta^* t_k}}\right)\epsilon} \frac{e^{-\tau} - 1}{\tau} d\tau \right] \quad (20b)$$

where $C = 0.577215$; Euler's constant.

In expressions (19b) and (20b) are equated to each other, i.e.,

$$\lim_{\epsilon \rightarrow 0} -\frac{1}{t_k} \ln \left[1 - \left(\frac{1 - e^{-\beta^* t_k}}{\epsilon} \right) \right] - \lim_{\epsilon \rightarrow 0} -\frac{1}{t_k} \ln \left[\left(\frac{\epsilon}{1 - e^{-\beta^* t_k}} \right) - 1 \right] =$$

$$\lim_{\epsilon \rightarrow 0} \frac{-1}{V_c} e^{\left(\frac{-\ln(V_k)}{(1-e^{-\beta^* t_k})}\right)\epsilon} \left[C + \ln \left[\frac{\ln(V_k)}{1 - e^{-\beta^* t_k}} \right] \epsilon + \int_0^{\left(\frac{\ln(V_k)}{1 - e^{-\beta^* t_k}}\right)\epsilon} \frac{e^{-\tau} - 1}{\tau} d\tau \right]$$

We expanding the a series and integrand in the value of β and the limit can be expressed then In the resulting equation retaining dominant $\ln(1/\epsilon)$ term only, after a little algebra, we get

$$\left(1 - \frac{t_k}{V_c} \right) \ln \frac{1}{\epsilon} = -\frac{t_k}{V_c} C - \frac{t_k}{V_c} \ln \left[\frac{\ln(V_k)}{1 - e^{-\beta^* t_k}} \right] - \ln[1 - e^{-\beta^* t_k}] \quad (21)$$

or, equivalently,

$$\epsilon = \exp \left\{ \ln[1 - e^{-\beta^* t_k}] - \frac{t_k}{V_c} \left(C + \ln \left(\frac{\ln(V_k)}{1 - e^{-\beta^* t_k}} \right) \right) \right\}$$

A simple substitution of ϵ into $-\frac{1}{t_k} \ln \left[1 - \left(\frac{1 - e^{-\beta^* t_k}}{\epsilon} \right) \right]$

yields the required asymptotic formula

$$\beta_k = -\frac{1}{t_k} \ln \left[1 - e^{\left(\frac{t_k}{V_c}\right)C / \left(\frac{t_k}{V_c}-1\right)} (\ln V^*(t))^{\left(\frac{t_k}{V_c}\right) / \left(\frac{t_k}{V_c}-1\right)} \right] \quad (22)$$



of (15) for large $V^*(t) = V_k$ with $\frac{t_k}{V_c \ln V^*(t)} = 1$
 Strangely, on account of (21), the above formula fails at $\frac{t_k}{V_c} = 1$. But when $\frac{t_k}{V_c} = 1$, we have

$$0 \leq \beta_k \leq \left(\frac{1 - e^{-\beta_k t_k}}{t_k \ln[V_k]} \right) \text{ from (16). Now setting } \left(\frac{1 - e^{-\beta_k t_k}}{\ln[V_k]} \right) = \frac{1}{\epsilon}$$

substitute this into equation (15), to get

$$\beta_k = \lim_{\epsilon \rightarrow 0} \frac{-1}{t_k} e^{-\epsilon} \int_{-\epsilon}^{\infty} \frac{e^{-z}}{z} dz$$

If we send ϵ to zero, on account of 8.212.1. in [45], we get

$$\beta_k = \lim_{\epsilon \rightarrow 0} -\frac{1}{t_k} \ln \left[1 - \left(\frac{\ln[V_k]}{\epsilon} \right) \right] \quad (23a)$$

$$= \lim_{\epsilon \rightarrow 0} -\frac{1}{t_k} \ln \left[1 - \left(\frac{\ln[V_k]}{\epsilon} \right) \right] - \lim_{\epsilon \rightarrow 0} -\frac{1}{t_k} \ln \left[\left(\frac{\epsilon}{\ln[V_k]} \right) - 1 \right] \quad (23b)$$

And also from equation (17),

$$\beta_k = \lim_{\epsilon \rightarrow 0} \frac{-1}{t_k} e^{-\epsilon} \int_{-\epsilon}^{\infty} \frac{e^{-z}}{z} dz \quad (24a)$$

$$= \lim_{\epsilon \rightarrow 0} \frac{-1}{t_k} e^{-\epsilon} \left[C + \ln(\epsilon) + \int_0^{\epsilon} \frac{e^{-\tau} - 1}{\tau} d\tau \right] \quad (24b)$$

In expressions (23b) and (24b) are equated to each other, i.e.,

$$\lim_{\epsilon \rightarrow 0} -\frac{1}{t_k} \ln \left[1 - \left(\frac{\ln[V_k]}{\epsilon} \right) \right] - \lim_{\epsilon \rightarrow 0} -\frac{1}{t_k} \ln \left[\left(\frac{\epsilon}{\ln[V_k]} \right) - 1 \right] =$$

$$= \lim_{\epsilon \rightarrow 0} \frac{-1}{t_k} e^{-\epsilon} \left[C + \ln(\epsilon) + \int_0^{\epsilon} \frac{e^{-\tau} - 1}{\tau} d\tau \right]$$

We expanding the a series and integrand in the value of β and the limit can be expressed. then in the resulting equation, as the dominant $\ln(1/\epsilon)$ term vanishes, we consider only ϵ -order term.

After a little algebra, we obtain

$$-\frac{1}{t_k} \ln[(\ln(V_k))] - \frac{\epsilon}{t_k \ln[V_k]} = \frac{1}{t_k} [(1 - \epsilon)(C - \epsilon)]$$

Or

$$\epsilon = \frac{\ln[(\ln(V_k))] + C}{C - 1 - (1/\ln(V_k))}$$

A simple substitution of ϵ into $-\frac{1}{t_k} \ln \left[1 - \left(\frac{\ln[V_k]}{\epsilon} \right) \right]$ yields the required asymptotic formula

$$\beta_k = -\frac{1}{t_k} \ln \left[1 - \frac{C \ln[V_k] - \ln[V_k] - 1}{\ln(\ln[V_k]) + C} \right], \quad (25)$$

of (15) for large $\ln[V_k]$ when $\frac{t_k}{V_c} = 1$. summing up, we

conclude the following.

THEOREM : For every fixed V_c, t_k and $V^*(t)$, let I_1 and I_2 be intervals defined by

$$I_1 = \left[0, -\frac{1}{t_k} \ln \left[1 - e^{\left(\frac{t_k}{V_c} \right)^C / \left(\frac{t_k}{V_c} - 1 \right)} \left(\ln V^*(t) \right)^{\left(\frac{t_k}{V_c} \right) / \left(\frac{t_k}{V_c} - 1 \right)} \right] \right]$$

$$I_2 = \left[0, -\frac{1}{t_k} \ln \left[1 - \frac{C \ln[V_k] - \ln[V_k] - 1}{\ln(\ln[V_k]) + C} \right] \right]$$

where $C = 0.577215$; Euler's constant. Suppose there exists a unique solution of (7) in I_1, I_2 respectively, when $\frac{t_m}{V_c}$

$\neq 1$ and $\frac{t_m}{V_c} = 1$ Then it is necessary that $\frac{t_m}{V_c \ln V^*(t)} > 1$.

Moreover, the asymptotic solution of (7) for a large $V^*(t) = V_k$ is given by (22), (25), respectively, when

$\frac{t_m}{V_c} \neq 1$ and $\frac{t_m}{V_c} = 1$.

Remark 1. Note that care must be taken while using the asymptotic solution (22) and (25),

since the condition when $t_k = t_m, V_k = e^{\frac{t_m}{V_c}}$ and when $t_k = V_c \ln(V_k), V_k = e^{\frac{t_m}{V_c}}$ and $t_k \neq t_m$ are crucially dependent on whether given $\frac{t_m}{V_c \ln V^*(t)} > 1$ or $\frac{t_m}{V_c \ln V^*(t)} < 1$.

Remark 2. From (4) it is easy to get the asymptotic formula of A using (22) and (25) for $V^*(t) = V_k$

Remark 3. We listed in table I the numerical values of asymptotic solution (22) for $V^*(t) = V_k$ for a comparison with that for a large $V^*(t)$

Remark 4. There were no recorded samples with $\frac{t_m}{V_c} = 1$ to compare with the asymptotic solution(25) for $V^*(t) = V_k$.

Remark 5. Asymptotic solution is useful in the study of qualitative behaviour of solution.

Remark 6. Using (13) or (14) from the values of V_c, t_m and $V^*(t)$ we can calculate unique β then using that β substitute in (4) we will get the value of A.

Remark 7. Using (22) or (25) we can find critical Gompertz parameter β_k for large number solid tumour growth cells, in the absence of specific volume data at particular time.

IV. CONCLUSION

The above theorem states that the asymptotic solution β of (7) is a continuous function in the variables $\frac{1}{V_c}, t_m$ and $V^*(t)$ from initial growth to critical growth time (or, $V^*(t) \geq V_k$). What happens to solution β when $V^*(t) \leq V_k$ (below the critical volume)? As we have already equated t_m with t_k when $\frac{t_m}{V_c \ln V^*(t)} > 1$, what then is the actual maximum life time(theoretically exists)

t^*_m ? Such $t^*_m > t_m$ exists, How to determine this t^*_m ? All these questions remain to be addressed.

The purpose of this discussion has been to address the issue of existence of critical Gompertz parameter β_k for large number solid tumour growth cells, in the absence of specific volume data at particular time. Such a method is necessary when attempting to estimate the growth deceleration rate parameter. From these analyses, we believe that our model and methods will provide a useful approach to prediction of experimental and clinical tumour growth. For further applications more research is needed.

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Table-1 Analysis of Theoretical Gompertz Functions in terms of t_m, t_k and $VRD(asympotic)$ and $VRD(critical)$ which depends on the tumour cell number at any time.

	tumour type	$\frac{1}{V_c}$	t_m days	t_k days	VRD (asympotic)	VRD (critical)	V_k
	Mouse:						
1.	Krebs	8.4567×10^{-6}	-	10.4429	0.15675 hours	0.135 hours	800×10^6 cells
2.	Ehrlich	8.0750×10^{-7}	-	331.1687	9.88071 hours	9.26 hours	1593×10^6 cells
3.	MC ₁ M, low dose	1.4616×10^{-6}	-	-	6.5103 hours	6.09 hours	467×10^6 cells
4.	6C ₃ HED, high dose	33.0076×10^{-5}	425.9917	170.2028	20.01456 hours	19.6 hours	890×10^6 cells
5.	6C ₃ HED, low dose	9.1051×10^{-5}	567.3520	141.5309	13.02839 hours	11.9 hours	4776×10^6 cells
6.	DBA lymphoma	4.8265×10^{-8}	276.5147	207.188	2.68454 hours	2.59 hours	1000×10^6 cells
7.	EI ₄ , low dose	6.0002×10^{-8}	438.213	317.4464	3.53419 hours	3.46 hours	1260×10^6 cells
8.	EI ₄ , low dose high dose	1.03297×10^{-5}	-	-	5.62582 hours	4.23 hours	1290×10^6 cells
9.	EO771	2.9392×10^{-4}	79.4302	32.9247	1.33441 days	1.08 days	31cm^3
10.	Osteosarcomas	12.729×10^{-5}	38.9177	18.2688	0.75516 days	0.719 days	$4:3 \text{cm}^3$
	Rat:						
11.	Walker, W26b1	5.3999×10^{-5}	341.2281	42.3284	4.03779 days	3.26 days	175 g
12.	Walker, W12a7	3.8778×10^{-5}	342.2281	51.0880	2.2779 days	2.07 days	212cm^3
13.	Walker, W10a6	2.3222×10^{-3}	59.1180	36.7349	2.02498 days	1.99 days	490cm^3
14.	Walker, W10b4	1.06732×10^{-4}	-	79.8275	6.09501 days	5.29 days	196cm^3
15.	R39Sarcoma, R3a7	1.5610×10^{-4}	42.4407	28.4854	0.56666 days	0.56 days	188cm^3
16.	7R39Sarcoma, R4c4	5.3331×10^{-3}	-	32.2799	1.32498 days	1.35 days	276cm^3
17.	R39Sarcoma, a7R3	1.6563×10^{-5}	-	62.7635	0.9851 days	0.97 days	202cm^3
18.	Flexne-Jobling	1.865×10^{-5}	-	43.9302	2.02099 days	1.84 days	18.3 g
	Rabbit:						
19.	Brown-Pearce	2.08294×10^{-3}	45.7709	28.9986	0.56168 days	0.576 days	$29:8 \text{cm}^3$

The source of data for each species is given in [9,42,46].

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